

# Far-Ultraviolet and Far-Infrared Bivariate Luminosity Function of Galaxies: Complex Relation between Stellar and Dust Emission

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Far-ultraviolet (FUV) and far-infrared (FIR) luminosity functions (LFs) of galaxies show a strong evolution from  $z = 0$  to  $z = 1$ , but the FIR LF evolves much stronger than the FUV one. The FUV is dominantly radiated from newly formed short-lived OB stars, while the FIR is emitted by dust grains heated by the FUV radiation field. It is known that dust is always associated with star formation activity. Thus, both FUV and FIR are tightly related to the star formation in galaxies, but in a very complicated manner. In order to disentangle the relation between FUV and FIR emissions, we estimate the UV-IR bivariate LF (BLF) of galaxies with *GALEX* and *AKARI* All-Sky Survey datasets. Recently we invented a new mathematical method to construct the BLF with given marginals and prescribed correlation coefficient. This method makes use of a tool from mathematical statistics, so called “copula”. The copula enables us to construct a bivariate distribution function from given marginal distributions with prescribed correlation and/or dependence structure. With this new formulation and FUV and FIR univariate LFs, we analyze various FUV and FIR data with *GALEX*, *Spitzer*, and *AKARI* to estimate the UV-IR BLF. The obtained BLFs naturally explain the nonlinear complicated relation between FUV and FIR emission from star-forming galaxies. Though the faint-end of the BLF was not well constrained for high- $z$  samples, the estimated linear correlation coefficient  $\rho$  was found to be very high, and is remarkably stable with redshifts (from 0.95 at  $z = 0$  to 0.85 at  $z = 1.0$ ). This implies the evolution of the UV-IR BLF is mainly due to the different evolution of the univariate LFs, and may not be controlled by the dependence structure.

**Key words:** Dust; galaxies: formation; galaxies: evolution; stars: formation; infrared; ultraviolet.

## 1. Introduction

Exploring the star formation history of galaxies is one of the most important topics in modern observational cosmology. Especially, the “true” absolute value of the cosmic star formation rate (hereafter SFR) has been of a central importance to understand the formation and evolution of galaxies.

However, it has been a difficult task for a long time because of dust extinction. Active star formation (SF) is always accompanied by dust production through various dust grain formation processes related to the final stage of stellar evolution (e.g., Dwek and Scalo(1980); Dwek 1998; Nozawa et al. 2003; Takeuchi et al. 2005c; Asano et al. 2012). Observationally, SFR of galaxies are, in principle, measured by the ultraviolet (UV) luminosity from massive stars because of their short lifetime ( $\sim 10^6$  yr) compared with the age of galaxies or the Universe. However, since the UV photons are easily scattered and absorbed by dust grains, SFR of galaxies is always inevitably affected by dust produced by their own SF activity. Since the absorbed energy is re-emitted at wavelengths of far-infrared (FIR), we need observations both at UV and FIR to have an unbiased view of their SF (e.g., Buat et al. 2005; Seibert et al. 2005; Cortese et al. 2006; Takeuchi et al. 2005a; Takeuchi et al. 2010; Haines et al. Bothwell et al. 2011; Hao et al. 2011).

After great effort of many researchers, the cosmic history of the SFR density is gradually converging at  $0 < z < 1$ . This “latter half” of the cosmic SFR is characterized by the rapid decline of the total SFR, especially the decrease of the contribution of dusty IR galaxies toward  $z = 0$ : While at  $z = 0$ , the contribution of the SFR hidden by dust is 50–60 %, it becomes  $> 90$  % at  $z = 1$  (Takeuchi et al. 2005a). This difference of decrease in SFR obtained from FUV and FIR has been already recognized in IR studies (e.g., Takeuchi et al. 2001a,b). Later works confirmed this dusty era of the Universe, and revealed that the dominance of the hidden SF continues even toward higher redshifts ( $z \sim 3$ ) (e.g., Murphy et al. 2011; Cucciati et al. 2012).

Then, a natural question arises: what does the different evolution at different wavelengths represent? To address this problem, it is very important to understand how we select sample galaxies and what we see in them. Each time we find some relation between different properties, we must understand clearly which is real (physical) and which is simply due to a selection effect. Some are detected at both bands, but some are detected only at one of the observed wavelength and appear as upper limits at the other wavelength. In previous studies it was often found that various claims were inconsistent with each other, mainly because they did not construct a well-controlled sample of FUV and FIR selected galaxies. Recently, thanks to new large surveys, some attempts to explore the SFR distribution of galaxies in bivariate way have been made, through the “SFR function” (e.g.,

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Buat et al. 2007, 2009, Haines et al. 2011; Bothwell et al. 2011). These works are based on the FUV and FIR LFs and their sum, but have not yet address their dependence on each other. To explore the bivariate properties of the SF in galaxies further, the proper tool is the UV-IR bivariate luminosity function (BLF).

However, constructing a BLF from two-wavelength data is not a trivial task. When we have a complete flux-limited<sup>1</sup> multiwavelength dataset, we can estimate a univariate luminosity function (LF) at each band, but what we want to know is *the dependence structure* between luminosities at different bands. Mathematically, this problem is rephrased as follows: can we (re)construct a multivariate probability density function (PDF) satisfying prescribed marginals? Although there is an infinite number of degrees of freedom to choose the original PDF, if we can model the dependence between variables, we can construct such a bivariate PDF. A statistical tool for this problem is the so-called “copula” (see Sec. 2 for the definition).

A copula has been extensively used in financial engineering, for instance, but until recently there were very few applications to astrophysical problems (e.g., Koen 2009; Benabed et al. 2009; Scherrer et al. 2010; Sato et al. 2010, 2011). Takeuchi (2010) introduced the copula to the estimation problem of a BLF. In this work, we apply the copula-based BLF estimation to the UV-IR two-wavelength dataset from  $z = 0$  to  $z = 1$ , using data from *IRAS*, *AKARI*<sup>2</sup>, *Spitzer*<sup>3</sup>, and *GALEX*<sup>4</sup>.

The layout of this paper is as follows. We define a copula, especially the Gaussian copula, and formulate the copula-based BLF in Sec. 2. We then describe our UV and IR data in Sec. 3. In Sec. 4, we first formulate the likelihood function for the BLF estimation. Then we show the results, and discuss the possible interpretation of the evolution of the UV-IR BLF. Section 5 is devoted to summary.

Throughout the paper we will assume  $\Omega_{M0} = 0.3$ ,  $\Omega_{\Lambda0} = 0.7$  and  $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ . The luminosities are defined as  $\nu L_\nu$  and expressed in solar units assuming  $L_\odot = 3.83 \times 10^{33} \text{ erg s}^{-1}$ .

## 2. The Bivariate Luminosity Function Based on the Copula

### 2.1 Copula: general definition

First, we briefly introduce a copula. A copula  $C(u_1, u_2)$  is defined as follows:

$$G(x_1, x_2) = C[F_1(x_1), F_2(x_2)] \quad (1)$$

where  $F_1(x_1)$  and  $F_2(x_2)$  are two univariate marginal cumulative distribution functions (DFs) and  $G(x_1, x_2)$  is a bivariate DF. We note that all bivariate DFs have this form and we can safely apply this method to any kind of bivariate DF estimation problem (Takeuchi 2010). In various applications,

we usually know the marginal DFs (or equivalently, PDFs) from the data. Then, the problem reduces to a statistical estimation of a set of parameters to determine the shape of a copula  $C(u_1, u_2)$ . In the form of the PDF,

$$\begin{aligned} g(x_1, x_2) &= \frac{\partial^2 C[F_1(x_1), F_2(x_2)]}{\partial x_1 \partial x_2} f_1(x_1) f_2(x_2) \\ &\equiv c[F_1(x_1), F_2(x_2)] f_1(x_1) f_2(x_2), \end{aligned} \quad (2)$$

where  $f_1(x_1)$  and  $f_2(x_2)$  are the PDFs of  $F_1(x_1)$  and  $F_2(x_2)$ , respectively. On the second line,  $c(u_1, u_2)$  is referred to as a differential copula.

### 2.2 Gaussian copula

Since the choice of copula is literally unlimited, we have to introduce a guidance principle. In many data analyses in physics, the most familiar measure of dependence might be the linear correlation coefficient  $\rho$ . Mathematically speaking,  $\rho$  depends not only on the dependence of two variables but also the marginal distributions, which is not an ideal property as a dependence measure. Even so, a copula having an explicit dependence on  $\rho$  would be convenient. In this work, we use a copula with this property, the Gaussian copula.

One of the natural candidate with  $\rho$  may be a copula related to a bivariate Gaussian DF (for other possibilities, see Takeuchi 2010). The Gaussian copula has an explicit dependence on a linear correlation coefficient by its construction. Let

$$\psi_1(x) = \frac{1}{\sqrt{2\pi}} \exp\left(-\frac{x^2}{2}\right), \quad (3)$$

$$\Psi_1 = \int_{-\infty}^x \psi_1(x') dx', \quad (4)$$

$$\psi_2(x_1, x_2; \rho) = \frac{1}{\sqrt{(2\pi)^2 \det \Sigma}} \exp\left(-\frac{1}{2} \mathbf{x}^T \Sigma^{-1} \mathbf{x}\right), \quad (5)$$

and

$$\Psi_2(x_1, x_2; \rho) = \int_{-\infty}^{x_1} \int_{-\infty}^{x_2} \psi_2(x'_1, x'_2) dx'_1 dx'_2, \quad (6)$$

where  $\mathbf{x} \equiv (x_1, x_2)^T$ ,  $\Sigma$  is a covariance matrix

$$\Sigma \equiv \begin{pmatrix} 1 & \rho \\ \rho & 1 \end{pmatrix}, \quad (7)$$

and superscript  $T$  stands for the transpose of a matrix or vector.

Then, we define a Gaussian copula  $C^G(u_1, u_2; \rho)$  as

$$C^G(u_1, u_2; \rho) = \Psi_2[\Psi_1^{-1}(u_1), \Psi_1^{-1}(u_2); \rho]. \quad (8)$$

<sup>1</sup> Or any other observational condition.

<sup>2</sup> URL: <http://www.ir.isas.ac.jp/ASTRO-F/index-e.html>.

<sup>3</sup> URL: <http://www.spitzer.caltech.edu/>.

<sup>4</sup> URL: <http://www.galex.caltech.edu/>.

The density of  $C^G$ ,  $c^G$ , is obtained as

$$\begin{aligned} c^G(u_1, u_2; \rho) &= \frac{\partial^2 C^G(u_1, u_2; \rho)}{\partial u_1 \partial u_2} \\ &= \frac{\partial^2 \Psi_2 [\Psi_1^{-1}(u_1), \Psi_1^{-1}(u_2); \rho]}{\partial u_1 \partial u_2} \\ &= \frac{\psi_2(x_1, x_2; \rho)}{\psi_1(x_1)\psi_1(x_2)} \\ &= \frac{1}{\sqrt{\det \Sigma}} \exp \left\{ -\frac{1}{2} [\Psi^{-1 T} (\Sigma^{-1} - \mathbf{I}) \Psi^{-1}] \right\}, \end{aligned} \quad (9)$$

where  $\Psi^{-1} \equiv [\Psi^{-1}(u_1), \Psi^{-1}(u_2)]^T$  and  $\mathbf{I}$  stands for the identity matrix.

### 2.3 Construction of the UV-IR BLF

In this work, we define the luminosity at a certain wavelength band by  $L \equiv \nu L_\nu$  ( $\nu$  is the corresponding frequency). Then the LF is defined as the number density of galaxies whose luminosity lies between a logarithmic interval  $[\log L, \log L + d \log L]$ :

$$\phi^{(1)}(L) \equiv \frac{dn}{d \log L}, \quad (10)$$

where we denote  $\log x \equiv \log_{10} x$  and  $\ln x \equiv \log_e x$ . For mathematical simplicity, we define the LF as being normalized, i.e.,

$$\int \phi^{(1)}(L) d \log L = 1. \quad (11)$$

Hence, this corresponds to a PDF. We also define a cumulative LF as

$$\Phi^{(1)}(L) \equiv \int_{\log L_{\min}}^{\log L} \phi^{(1)}(L') d \log L', \quad (12)$$

where  $L_{\min}$  is the minimum luminosity of galaxies considered. This corresponds to the DF. If we denote the univariate LFs as  $\phi_1^{(1)}(L_1)$  and  $\phi_2^{(1)}(L_2)$ , then the BLF  $\phi^{(2)}(L_1, L_2)$  is described by a differential copula  $c(u_1, u_2)$  as

$$\phi^{(2)}(L_1, L_2) = c [\phi_1^{(1)}(L_1), \phi_2^{(1)}(L_2)]. \quad (13)$$

For the Gaussian copula, the BLF is obtained as

$$\begin{aligned} \phi^{(2)}(L_1, L_2; \rho) &= \frac{1}{\sqrt{\det \Sigma}} \exp \left\{ -\frac{1}{2} [\Psi^{-1 T} (\Sigma^{-1} - \mathbf{I}) \Psi^{-1}] \right\} \\ &\quad \times \phi_1^{(1)}(L_1) \phi_2^{(1)}(L_2), \end{aligned} \quad (14)$$

where

$$\Psi^{-1} = [\Psi^{-1}(\Phi_1^{(1)}(L_1)), \Psi^{-1}(\Phi_2^{(1)}(L_2))]^T. \quad (15)$$

For the UV LF, we adopt the Schechter function (Schechter 1976).

$$\phi_1^{(1)}(L) = (\ln 10) \phi_{*1} \left( \frac{L}{L_{*1}} \right)^{1-\alpha_1} \exp \left[ -\left( \frac{L}{L_{*1}} \right) \right]. \quad (16)$$

For the IR, we use the analytic form for the LF proposed by Saunders et al. (1990),

$$\begin{aligned} \phi_2^{(1)}(L) &= \phi_{*2} \left( \frac{L}{L_{*2}} \right)^{1-\alpha_2} \exp \left\{ -\frac{1}{2\sigma^2} \left[ \log \left( 1 + \frac{L}{L_{*2}} \right) \right]^2 \right\}. \end{aligned} \quad (17)$$

We use the re-normalized version of eqs. (17) and (16) so that they can be regarded as PDFs, as mentioned above.

### 2.4 Selection effects: another benefit of a copula BLF

Another benefit of copula is that it is easy to incorporate observational selection effects which always exist in any kind of astronomical data. In a bi(multi)variate analysis, there are two categories of observational selection effects.

#### 1) Truncation

We do not know if a source would exist below a detection limit.

#### 2) Censoring

We know there is a source, but we have only an upper (sometimes lower) limit for a certain observable.

As we mentioned above, we have to deal with both of these selection effects carefully to construct a BLF from data at the same time. It is terribly difficult to incorporate these effects by heuristic methods, particularly for a nonparametric methods (Takeuchi 2012, in preparation). In contrast, since we have an analytic form for the BLF, the treatment of upper limits is much more straightforward (Takeuchi 2010). We show how it is treated in the likelihood function in Sec. 4.

## 3. Data

### 3.1 UV-IR data construction

We have constructed a dataset of galaxies selected at FUV and FIR by *GALEX* and *IRAS* for  $z = 0$ . At higher redshifts, *GALEX* and *Spitzer* data are used for  $z = 0.7$  and EIS (ESO Imaging Survey) and *Spitzer* for  $z = 1.0$  in the Chandra Deep Field South (CDFs). We explain the details of each sample in the following.

In the Local Universe, we used the sample compiled by Buat et al. (2007). This sample was constructed based on *IRAS* all-sky survey and *GALEX* All-Sky Imaging Survey (AIS). This dataset consists of UV- and IR-selected samples. The UV-selection was made by the *GALEX* FUV ( $\lambda = 1530 \text{ \AA}$ ) band with  $FUV < 17 \text{ mag}$  (hereafter, all magnitudes are presented in AB mag). This corresponds to the luminosity lower limit of  $L_{FUV} > 10^8 L_\odot$ . Redshift information was taken from HyperLEDA (Paturel et al. 2003) and NED. The IR-selection is based on the *IRAS* PSCz (Saunders et al. 2000). The detection limit of the PSCz is  $S_{60} > 0.6 \text{ Jy}$ . Redshifts of all PSCz galaxies were measured, and most of their redshifts are  $z < 0.05$ . All UV fluxes were remeasured with the package we have developed (Iglesias-Páramo et al. 2006) to avoid the shredding of galaxies. Details of the sample construction are explained in Buat et al. (2007). The number of the UV- and IR-selected samples are 606 and 644, respectively.

We also constructed a new, much larger sample of IR-selected galaxies by *AKARIFIS* All-Sky Survey. We started

from the *AKARI* FIS bright source catalog (BSC) v.1 from the *AKARI* all sky survey (Yamamura et al. 2010). This sample is selected at *WIDE-S* band ( $\lambda = 90 \mu\text{m}$ ) of the *AKARI* FIS (Kawada et al. 2007). The detection limit is  $S_{90} > 0.2 \text{ Jy}$ . We first selected *AKARI* sources in the SDSS footprints to have the same solid angle with a forthcoming corresponding UV-selected sample we are preparing with *GALEX*-SDSS.<sup>5</sup> Then, to have a secure sample of galaxies with redshift data, we made a cross match of *AKARI* sources with the Imperial *IRAS*-FSC Redshift Catalogue (IIFSCz), a redshift catalog recently published (Wang and Rowan-Robinson 2009), with a search radius of 36 arc-sec. Since about 90 % of galaxies in the IIFSCz have spectroscopic or photometric redshifts at  $S_{60} > 0.36 \text{ Jy}$ , the depth of the sample is defined by this matching. This determines the redshift range of this dataset, to be approximately  $z < 0.1$ . We measured the FUV and NUV flux densities with the same procedure as the *IRAS*-*GALEX* sample. The detection limits at FUV and NUV of this sample are 19.9 mag and 20.8 mag (Morrissey et al. 2007). A corresponding UV-selected sample is under construction, hence we have only IR-selected sample. The number of galaxies is 3,891. For more information on this sample and properties of galaxies, see Sakurai et al. (2012).

At higher- $z$ , our samples are selected in the CDFS. *GALEX* observed this field at FUV and NUV (2300 Å) as a part of its deep imaging survey. We restricted the field to the subfield observed by *Spitzer*/MIPS as a part of the GOODS key program (Elbaz et al. 2007) to have the corresponding IR data. The area of the region is  $0.068 \text{ deg}^2$ . Precise description of our high- $z$  samples are presented in Buat et al. (2009) and Burgarella et al. (2006).

At  $z = 0.7$ , the NUV-band corresponds to the rest-frame FUV of the sample galaxies. We thus constructed the sample based on NUV-selection. Redshifts were taken from the COMBO-17 survey (Wolf et al. 2004), and we have defined the  $z = 0.7$  sample as those with redshifts of  $0.6 < z < 0.8$ . Data reduction and redshift association are explained in Burgarella et al. (2006). We truncated the sample at NUV = 25.3 mag so that more than 90 % of the *GALEX* sources are identified in COMBO-17 with  $R < 24 \text{ mag}$ . We set the MIPS 24  $\mu\text{m}$  upper limit as 0.025 mJy for undetected sources. For the IR-selected sample, we based on the GOODS *Spitzer*/MIPS 24  $\mu\text{m}$  sample and matched the *GALEX* and COMBO-17 sources. The sizes of the UV- and IR-selected samples are 340 and 470, respectively.

In contrast to  $z = 0.7$ , since NUV of *GALEX* corresponds to 1155 Å in the rest frame of galaxies at  $z = 1.0$ , we cannot use NUV as the primary selection band as rest-frame FUV. Instead, we selected galaxies in the *U*-band from the EIS survey (Armouts et al. 2001) that covers the CDFS/GOODS field. We then cross-matched the *U*-band sources with the COMBO-17 sample to obtain redshifts. For the  $z = 1$  sample, the range of redshifts is  $0.8 < z < 1.2$ . We set the limit at  $U = 24.3 \text{ mag}$  to avoid source confusion. The IR

flux densities were taken from *Spitzer* MIPS 24  $\mu\text{m}$  data, and the same upper limit as the  $z = 0.7$  sample is put for non-detections. Again the IR-selected sample was constructed from the GOODS and matched the *GALEX* and COMBO-17 sources, as the  $z = 0.7$  sample. The sizes of the  $z = 1.0$  UV and IR-selected samples are 319 and 1033, respectively.

The characteristics of these samples are summarized in Table 1.

### 3.2 Far-UV and total IR luminosities

We are interested in the SF activity of galaxies and its evolution. Hence, luminosities of galaxies representative of SF activity would be ideal. As for the directly visible SF, obviously the UV emission is appropriate for this purpose. We define the FUV luminosity of galaxies  $L_{\text{FUV}}$  as  $L_{\text{FUV}} \equiv \nu L_{\nu} @ \text{FUV}$ . For  $z = 0$  galaxies, FUV corresponds to 1530 Å. At higher redshifts, as we explained,  $L_{\text{FUV}}$  is calculated from the NUV flux density at  $z = 0.7$  and *U*-band flux density at  $z = 1.0$ , respectively.

In contrast, at IR, the luminosity related to the SF activity is the one integrated over a whole range of IR wavelengths ( $\lambda = 8\text{--}1000 \mu\text{m}$ ),  $L_{\text{TIR}}$ , where the subscript TIR stands for the total IR. For the Local *IRAS*-*GALEX* sample, it is quite straightforward to define  $L_{\text{TIR}}$  since the *IRAS* galaxies are selected at 60  $\mu\text{m}$ . We adopted a formula  $L_{\text{TIR}} = 2.5\nu L_{\nu} @ 60 \mu\text{m}$ . This rough approximation is, in fact, justified by spectral energy distributions (SEDs) of galaxies with *ISO* 160  $\mu\text{m}$  observations (Takeuchi et al. 2006).

Also, for the *AKARI*-*GALEX* sample, we have excellent flux density data at FIR. We used the following TIR estimation formula from *AKARI* two wide bands (Takeuchi et al. 2005b, 2010),

$$L_{\text{TIR}} = \Delta\nu(\text{WIDE-S})L_{\nu}(90 \mu\text{m}) + \Delta\nu(\text{WIDE-L})L_{\nu}(140 \mu\text{m}), \quad (18)$$

where

$$\begin{aligned} \Delta\nu(\text{WIDE-S}) &= 1.47 \times 10^{12} [\text{Hz}] \\ \Delta\nu(\text{WIDE-L}) &= 0.831 \times 10^{12} [\text{Hz}]. \end{aligned} \quad (19)$$

However, since deep FIR data of higher- $z$  galaxies are not easily available to date, we should rely on the data taken by *Spitzer* MIPS 24  $\mu\text{m}$ . we use conversion formulae from MIR luminosity to  $L_{\text{TIR}}$ :

$$\log L_{\text{TIR}}[L_{\odot}] = 1.23 + 0.972 \log L_{15}[L_{\odot}], \quad (20)$$

$$\begin{aligned} \log L_{\text{TIR}}[L_{\odot}] &= 2.27 + 0.707 \log L_{12}[L_{\odot}] \\ &\quad + 0.0140 (\log L_{12}[L_{\odot}])^2, \end{aligned} \quad (21)$$

which are an updated version of the formulae proposed by Takeuchi et al. (2005b) and also used by Buat et al. (2009). Here  $L_{15}$  and  $L_{12}$  are luminosities  $\nu L_{\nu} @ 15 \mu\text{m}$  and  $12 \mu\text{m}$ , i.e.,  $24/(1+z)$  at  $z = 0.7$  and  $1.0$ , respectively. The estimated  $L_{\text{TIR}}$  slightly depends on which kind of conversion formula is used, but for our current purpose, it does not affect our conclusion and we do not discuss extensively here. Intercomparison of the MIR-TIR conversion formulae can be found in Buat et al. (2009).

<sup>5</sup> This step is not a mandatory in this study, but we are planning to make an extension of this analysis with the UV-selected data, and this step will make the analysis easier with respect to the treatment of survey areas when the UV-selected data will be ready.

Table 1. Sample description

Redshift	[range]	Primary selection	Number
$z = 0$	[0, 0.05]	UV ( <i>GALEX</i> FUV)	606
		IR ( <i>IRAS</i> 60 $\mu\text{m}$ )	644
	[0, 0.1]	IR ( <i>AKARI WIDE-S</i> )	3891
$z = 0.7$	[0.6, 0.8]	UV ( <i>GALEX</i> NUV)	340
		IR ( <i>Spitzer</i> MIPS 24 $\mu\text{m}$ )	470
$z = 1.0$	[0.8, 1.2]	UV (EIS <i>u</i> -band)	319
		IR ( <i>Spitzer</i> MIPS 24 $\mu\text{m}$ )	1033

This situation will be greatly improved with *Herschel*<sup>6</sup> data. We will leave this as our future work with *Herschel* H-ATLAS (Eales et al. 2010) and GAMA (Galaxy And Mass Assembly: Driver et al. 2011) data (Takeuchi et al. 2012, in preparation).

## 4. Results and Discussion

### 4.1 FUV and TIR univariate LFs

In order to estimate the UV-IR BLF, first we have to examine our setting for the FUV and TIR univariate LFs. The *IRAS-GALEX* sample, the validity of the Local univariate LFs are already proved [see Fig. 3 of Buat et al. (2007)]. Hence, we can safely use the analytic formulae of FUV and TIR LFs at  $z = 0$  [eqs. (16) and (17)].

We use the Schechter parameters presented by Wyder et al. (2005) for *GALEX* FUV:  $(\alpha_1, L_{*1}, \phi_{*1}) = (1.21, 1.81 \times 10^9 h^{-2} [L_\odot], 1.35 \times 10^{-2} h^3 [\text{Mpc}^{-3}])$ . For the TIR, we used the parameters  $(\alpha_2, L_{*2}, \phi_{*2}, \sigma) = (1.23, 4.34 \times 10^8 h^{-2} [L_\odot], 2.34 \times 10^{-2} h^3 [\text{Mpc}^{-3}], 0.724)$  (Takeuchi et al. 2003) obtained from the *IRAS* PSC $z$  galaxies (Saunders et al. 2000), and multiplied  $L_{*1}$  with 2.5 to convert the 60- $\mu\text{m}$  LF to the TIR LF.

For higher redshifts, ideally we should estimate the univariate FUV and TIR LFs simultaneously with the BLF estimation. It is, however, quite difficult for our current samples because of the limited number of galaxies. We instead used the LF parameters at  $z = 0.7$  and  $1.0$  obtained by previous studies on univariate LFs and modeled the FUV and TIR univariate LFs at these redshifts and examined their validity with nonparametric LFs estimated from the data. We use the parameters compiled by Takeuchi et al. (2005a). Parameters of the evolution of the TIR LF are obtained by approximating the evolution in the form

$$\phi_2^{(1)}(L_2, z) = \mathfrak{g}(z)\phi_{2,0}^{(1)}\left[\frac{L_2}{\mathfrak{f}(z)}\right] \quad (22)$$

where  $\phi_{2,0}^{(1)}(z)$  is the local functional form of the TIR LF. Le Floc'h et al. (2005) assumed a power-law form for the evolution functions as

$$\mathfrak{f}(z) = (1+z)^\Omega, \quad \mathfrak{g}(z) = (1+z)^\mathfrak{P} \quad (23)$$

and obtained  $\mathfrak{P} = 0.7$  and  $\Omega = 3.2$ , with  $\alpha$  remaining constant. The Schechter parameters at  $z = 0.7$  and  $1.0$  for

the FUV LF are directly estimated by Arnouts et al. (2005) and we adopt their values (see Table 1 of Takeuchi et al. 2005a).

Then we estimated the FUV and TIR LFs with the step-wise maximum likelihood method and the variant of  $C^-$  method from our CDFS multiwavelength data (for the estimation method, see, e.g., Takeuchi et al. 2000; Johnston 2011, and references therein). The obtained univariate LFs at  $z = 0.7$  and  $1.0$  are presented in Figs. 1 and 2. We also show the analytic model LFs in these figures. We note that a well-known large density enhancement locates in the CDFS at  $z = 0.7$  (e.g., Salimbeni et al. 2009), and thus we have renormalized the LFs to remove the effect of the overdensity. In the following analysis, we normalize the univariate LFs according to eq. (11) so that we can treat the univariate LFs as PDFs, hence this does not affect the following analysis at all.

Because of the small sample size, the LF shape does not perfectly agree with the supposed functional forms, but the nonparametric LFs are acceptably similar to the analytic functions both for the FUV and TIR at each redshift. We stress that the analytic functions are *not* the fit to the data, but estimated from other studies. This implies that the estimated evolutionary parameters of the LFs work generally well. Thus, we can use the higher- $z$  univariate LFs as the marginal PDFs for the estimation of the UV-IR BLFs.

### 4.2 Copula likelihood for the BLF estimation

By using the estimated univariate FUV and TIR LFs as given marginals, we can estimate only one parameter, the linear correlation  $\rho$  by maximizing the likelihood function. The structure of a set of two-band selected data is  $(L_{\text{FUV}}^{i_k}, j_{\text{UV}}^{i_k}, L_{\text{TIR}}^{i_k}, j_{\text{IR}}^{i_k}), i_k = 1 \dots n_k$ . Here,  $j_{\text{band}}$  (band: UV or IR) stands for the upper limit flag:  $j_{\text{band}} = 0$ : detection and  $j_{\text{band}} = -1$ : upper limit. Another index  $k$  is the indicator of the selected band, i.e.,  $k = 1$  means a sample galaxy is selected at UV and  $k = -1$  means it is selected at

<sup>6</sup> URL: <http://herchel.esac.esa.int/>.

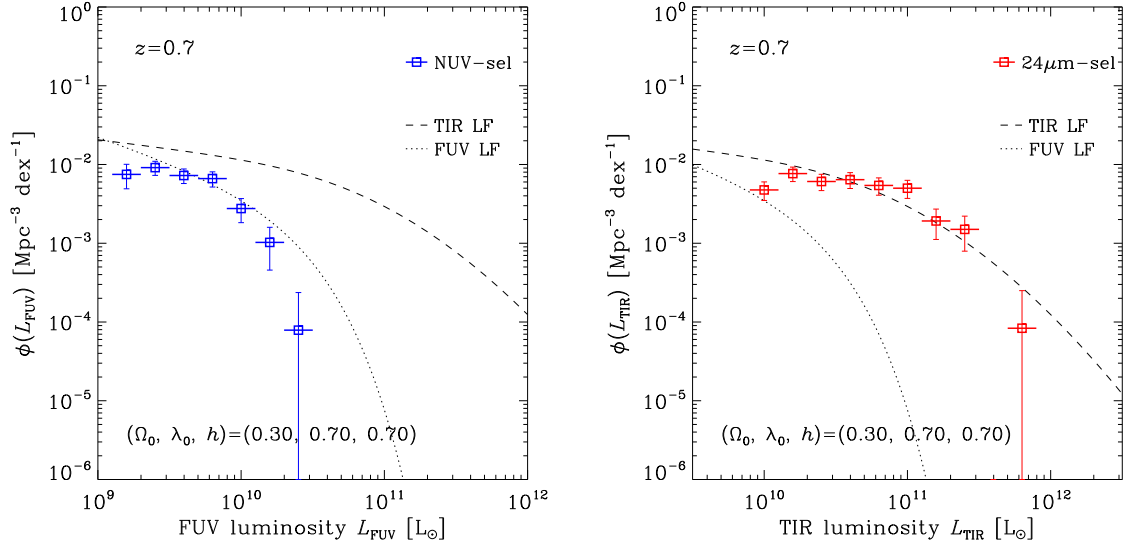


Fig. 1. The FUV and TIR LFs at  $z = 0.7$  (left and right panel, respectively). Open squares are a nonparametric LFs estimated from our CDFS multiwavelength data with *GALEX* NUV- and *Spitzer* MIPS 24  $\mu\text{m}$ -selections. Dotted lines represent the FUV LFs at this redshift bin taken from Arnouts et al. (2005). Dashed lines depict the TIR LFs derived from the evolutionary parameters at  $z = 0.7$  given by Le Floch et al. (2005). Because of a well-known large density enhancement at this redshift, we renormalized the LF to remove the effect of the overdensity.

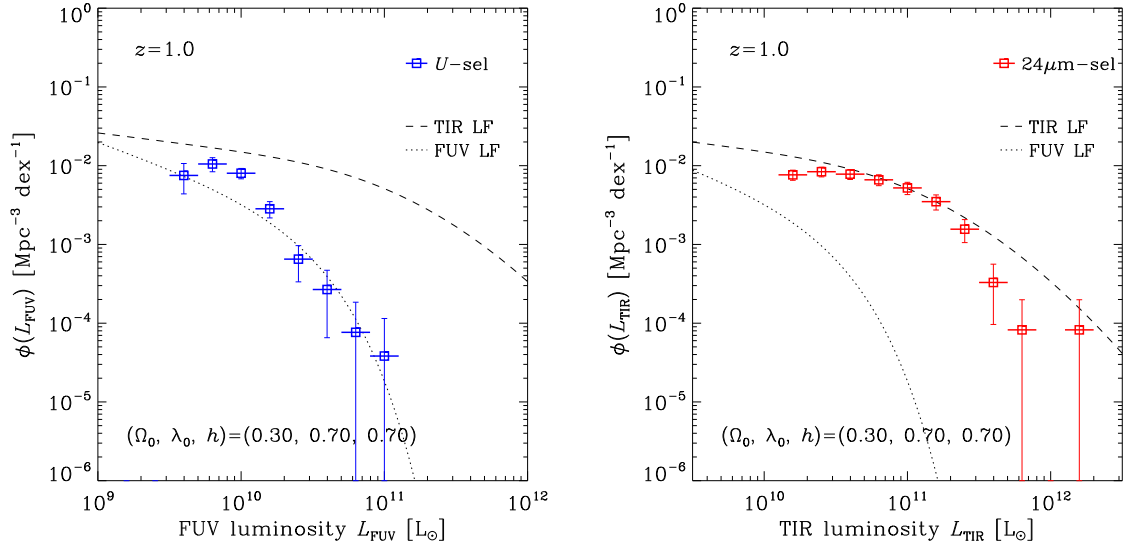


Fig. 2. The FUV and TIR LFs at  $z = 1.0$ . Symbols are the same as in Fig. 1, but FUV samples are selected at *U*-band.

IR. The likelihood function is as follows.

$$\begin{aligned}
 & \ln \mathcal{L}(L_{\text{FUV}}^{i_k}, L_{\text{TIR}}^{i_k} | i_k = 1, \dots, n_k, k = 1, -1) \\
 &= \sum_{k=\begin{smallmatrix} 1 : \text{UV sel} \\ -1 : \text{IR sel} \end{smallmatrix}} \sum_{i_k=1}^{n_k} \left\{ \ln [p^{\text{det}}(L_{\text{FUV}}^{i_k}, L_{\text{TIR}}^{i_k})]^{(1+j_{\text{UV}}^{i_k})(1+j_{\text{IR}}^{i_k})} \right. \\
 & \quad + \ln [p^{\text{UL:UV}}(L_{\text{FUV}}^{i_k}, L_{\text{TIR}}^{i_k})]^{\frac{(1+k)(-j_{\text{UV}}^{i_k})}{2}} \\
 & \quad \left. + \ln [p^{\text{UL:IR}}(L_{\text{FUV}}^{i_k}, L_{\text{TIR}}^{i_k})]^{\frac{(1-k)(-j_{\text{IR}}^{i_k})}{2}} \right\},
 \end{aligned}
 \tag{24}$$

where  $p^{\text{det}}(L_{\text{FUV}}^{i_k}, L_{\text{TIR}}^{i_k})$  is the probability for  $i_k$ th galaxy to be detected at both bands and to have luminosities  $L_{\text{FUV}}^{i_k}$  and  $L_{\text{TIR}}^{i_k}$ ,

$$p^{\text{det}}(L_{\text{FUV}}^{i_k}, L_{\text{TIR}}^{i_k}) \equiv \frac{\phi^{(2)}(L_{\text{FUV}}^{i_k}, L_{\text{TIR}}^{i_k}; \rho)}{\int_{L_{\text{FUV}}^{\text{lim}}(z_{i_k})}^{\infty} \int_{L_{\text{TIR}}^{\text{lim}}(z_{i_k})}^{\infty} \phi^{(2)}(L'_{\text{FUV}}, L'_{\text{TIR}}; \rho) dL'_{\text{TIR}} dL'_{\text{FUV}}},
 \tag{25}$$

$p^{\text{UL:UV}}(L_{\text{FUV}}^{i_k}, L_{\text{TIR}}^{i_k})$  is the probability for  $i_k$ th galaxy to be detected at IR band and have a luminosity  $L_{\text{TIR}}^{i_k}$ , but not detected at UV band and only an upper limit  $L_{\text{FUV},j_k}^{i_k}$  is

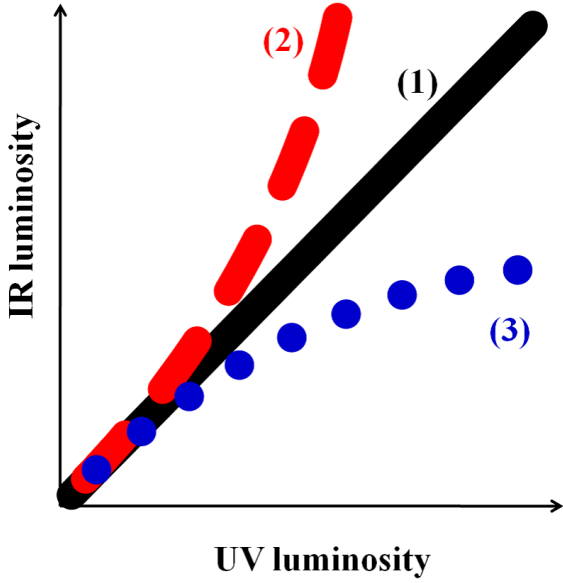


Fig. 3. Schematic BLF. (1) Diagonal: The energy from SF is emitted equally at UV and IR with any SF activity. (2) Upward: The more active the SF in a galaxy is, the more luminous at the IR (dusty SF). (3) Downward: The more active the SF is, the more luminous at the UV (transparent SF).

available,

$$p^{\text{UL: UV}}(L_{\text{FUV}}^{i_k}, L_{\text{TIR}}^{i_k}) \equiv \frac{\int_0^{L_{\text{FUV}}^{i_k, j_k}} \phi^{(2)}(L'_{\text{FUV}}, L_{\text{TIR}}^{i_k}) dL'_{\text{FUV}}}{\int_0^\infty \int_{L_{\text{TIR}}^{\text{lim}}(z_{i_k})}^\infty \phi^{(2)}(L'_{\text{FUV}}, L'_{\text{TIR}}; \rho) dL'_{\text{TIR}} dL'_{\text{FUV}}}, \quad (26)$$

and  $p^{\text{UL: IR}}(L_{\text{FUV}}^{i_k}, L_{\text{TIR}}^{i_k})$  is the probability for  $i_k$ th galaxy to be detected at UV band and have a luminosity  $L_{\text{FUV}}^{i_k}$ , but not detected at IR band and only an upper limit  $L_{\text{TIR}, j_k}^{i_k}$  is available,

$$p^{\text{UL: IR}}(L_{\text{FUV}}^{i_k}, L_{\text{TIR}}^{i_k}) \equiv \frac{\int_0^{L_{\text{TIR}, j_k}^{i_k}} \phi^{(2)}(L_{\text{FUV}}^{i_k}, L'_{\text{TIR}}) dL'_{\text{TIR}}}{\int_{L_{\text{FUV}}^{\text{lim}}(z_{i_k})}^\infty \int_0^\infty \phi^{(2)}(L'_{\text{FUV}}, L'_{\text{TIR}}; \rho) dL'_{\text{TIR}} dL'_{\text{FUV}}}. \quad (27)$$

The denominator in eq. (25) is introduced to take into account the truncation in the data by observational flux selection limits at both bands (e.g., Sandage et al. 1979; Johnston 2011). We should also note that it often happens that the same galaxies are included in both the UV- and IR-selected sample. In such a case, we should count the same galaxies only once to avoid double counting of them. Practically, such galaxies are included in any of the samples, because they are detected at both bands and are symmetric between UV- and IR-selections.

In the estimation procedure, we estimated the univariate LFs and their evolutions first, and then used these parameters

to estimate the BLF and its evolution. One might wonder if this compromising method would introduce some bias in the estimation. Since we fix the form of the maximum likelihood estimator eq. (24), this two-step estimation instead of simultaneous estimation does not bias the result, especially the dependence structure between UV and IR, unless the assumed univariate LF shape would be significantly different from the real one. We have already seen that the stepwise maximum likelihood estimation gave nonparametric LFs which reasonably agree with the assumed Schechter or Saunders function. Then, we can safely rely on the result for further discussion.

However, we should note that the two-step estimation gives conditional errors of the parameters only for each step, not the marginal ones. Then, the errors for each parameters are significantly underestimated. If we want to discuss the error values more precisely, we need a larger dataset and must use the simultaneous estimation method. This will be done in the future works.

#### 4.3 The BLF and its evolution

Using the Gaussian copula, now we can estimate the BLF. The visible and hidden SFRs should be directly reflected to this function. Dust is produced by SF activity, but also destroyed by SN blast waves as a result of the SF. Many physical processes are related to the evolution of the dust amount. Thus, first of all, we should describe statistically how it evolved, as stated in Introduction.

First, we summarize how to interpret the UV-IR BLF schematically in Fig. 3. First, we see the case that the ridge of the BLF is straight and diagonal [see (1) in Fig. 3]. This means that the energy from SF is emitted equally at UV and IR with any SF activity. If the relation is diagonal but has an offset horizontally or vertically, this suggests that a constant fraction of energy is absorbed by dust and reemitted. Second, if the ridge is curved upward, it means that the more active the SF in a galaxy is, the more luminous at the IR [dusty SF: (2) in Fig. 3]. Third, if the ridge is bent downward, the more active the SF is, the more luminous at UV [‘transparent’ SF: (3) in Fig. 3].

Now we show the estimated BLFs in Figs. 4–7. In the Local Universe, the BLF is quite well constrained. The estimated correlation coefficient  $\rho$  is very high:  $\rho = 0.95 \pm 0.04$  for *IRAS-GALEX* and  $\rho = 0.95 \pm 0.006$  for *AKARI-GALEX* datasets. The apparent scatter of the  $L_{\text{FUV}}-L_{\text{TIR}}$  is found to be due to the nonlinear shape of the ridge of the BLF. This bent shape of the BLF was implied by preceding studies (Martin et al. 2005), and we could quantify this feature. The copula BLF naturally reproduced it. For the *AKARI-GALEX* sample, only the IR-selected sample is available at this moment. This, however, does not bias the BLF estimation since the information from censored data points are incorporated properly in eq. (24). This can be directly tested, for example, by using only one of the UV and IR-selected data for the  $z = 0$  BLF estimation with *IRAS-GALEX* dataset. Both one-band estimations with UV and IR-selected data yield  $\rho = 0.95 \pm 0.07$ . Obviously the error becomes larger because of fewer data, but the estimate itself remain unchanged.

At higher redshifts ( $z = 0.7-1.0$ ), the linear correlation remains tight ( $\rho \simeq 0.91 \pm 0.06$  at  $z = 0.7$  and  $\rho \simeq 0.86 \pm 0.05$  at  $z = 1.0$ ), even though it is difficult to constrain the



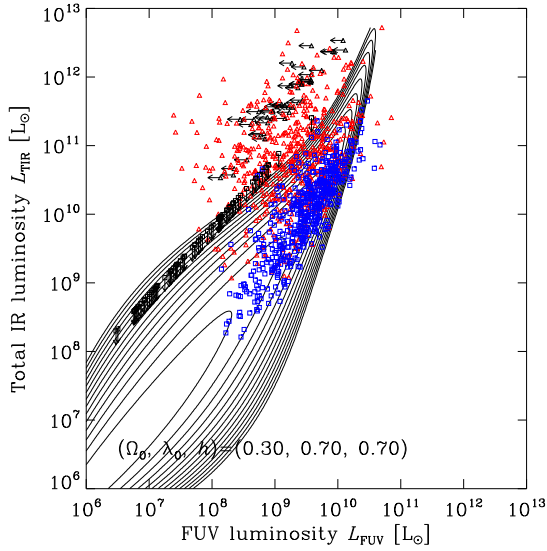


Fig. 4. The BLF of galaxies from *IRAS* and *GALEX* at  $z = 0$ . Contours are the analytic model constructed by a Gaussian copula and univariate FUV and TIR LFs. Open squares represent the UV-selected sample from *GALEX*, while open triangles are the IR-selected ones from *IRAS*. Squares and triangles with arrows mean that they are upper limits at FUV and FIR, respectively.

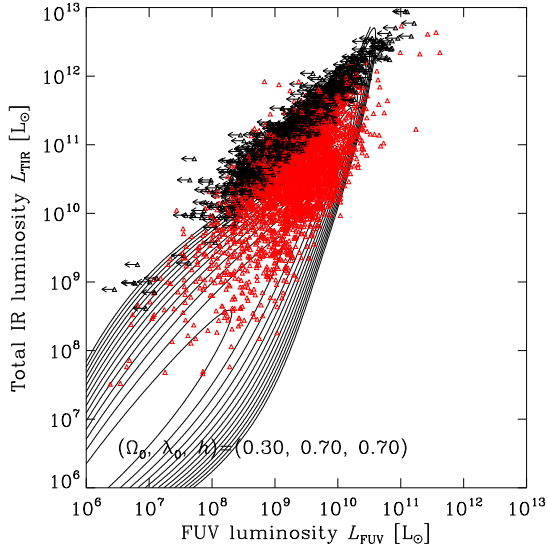


Fig. 5. The same as Fig. 4 but from *AKARI* and *GALEX*. Open triangles represent the IR-selected samples from *AKARI*. Since there is no UV-selected sample in this figure, we only show the IR-selected ones.

low-luminosity end from the data in this analysis (*Spitzer*-*GALEX* in the CDFS). The distribution of upper limits in Figs. 6 and 7 looks different from that in Fig. 4. Since we have restricted the redshifts of galaxies in these samples, the redshift restriction gives approximately constant luminosity limits both at FUV and FIR. This gives the “L-shaped” distribution of upper limits seen in these figures.

Though the whole shape cannot be perfectly determined by the current data, we find that  $\rho$  in the copula LF is high and remarkably stable with redshifts (from 0.95 at  $z = 0$  to

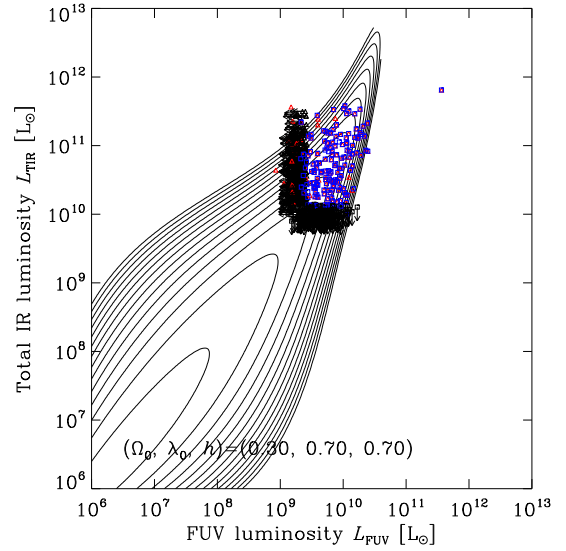


Fig. 6. The same as Fig. 4 but from *Spitzer* and *GALEX* at  $z = 0.7$ . Symbols are essentially the same as in Fig. 4, but IR-selected samples are from *Spitzer*/MIPS.

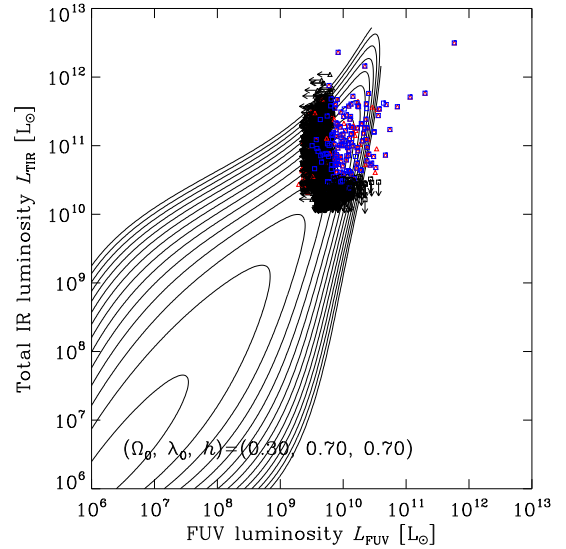


Fig. 7. The same as Fig. 4 but from *Spitzer* and *GALEX* at  $z = 1.0$ . Again symbols are essentially the same as in Fig. 4, but IR-selected samples are from *Spitzer*/MIPS.

0.85 at  $z = 1.0$ ). This implies the evolution of the UV-IR bivariate LF is mainly due to the different evolution of the univariate LFs, and may not be controlled by the evolution of the dependence structure.

Forthcoming better data in the future will improve various aspects of the BLF estimation. First, *Herschel*, ALMA, and *SPICA*<sup>7</sup> will provide us with direct estimations of  $L_{\text{TIR}}$  of galaxies especially at high redshifts, since they will detect these galaxies at longer wavelengths than  $24 \mu\text{m}$ . This allows us to avoid a possible bias of the estimated  $L_{\text{TIR}}$  caused by

<sup>7</sup> URL: [http://www.ir.isas.jaxa.jp/SPICA/SPICA\\_HP/index\\_English.html](http://www.ir.isas.jaxa.jp/SPICA/SPICA_HP/index_English.html).



the 24  $\mu\text{m}$ -based extrapolation of the SED.

Second, when we have deeper data at UV and IR, we will be able to extend the luminosity range toward much lower luminosities both at FUV and FIR. Then, the faint-end structure of the BLF will be much more tightly determined. Further, it is not only for the determination of the faint end, but deeper data will give some insights for the choice of the copula functional form. In this study, we did not try to examine if the Gaussian copula would be a proper choice as a model of the UV-IR BLF. At  $z = 0$ , when UV-selected data comparably large as the *AKARI* sample are ready, we will be able to specify (a family of) copulas appropriate for this analysis. At higher redshifts, our current data are not deep enough to have any claim on the faint-end dependence structure of the BLF, but again *Herschel*, *ALMA*, and *SPICA* data together with ground-based optical or *JWST* ones will enable us to examine which copula would be appropriate to describe the BLF, and constrain the evolution of the BLF along the whole history of galaxy evolution in the Universe.

## 5. Conclusion

To understand the visible and hidden SF history in the Universe, it is crucial to analyze multiwavelength data in a unified manner. The copula is an ideal tool to combine two marginal univariate LFs to construct a bivariate LFs. It is straightforward to extend this method to multivariate DFs.

- 1) The Gaussian copula LF is sensitive to the linear correlation parameter  $\rho$ .
- 2) Even so,  $\rho$  in the copula LF is remarkably stable with redshifts (from 0.95 at  $z = 0$  to 0.85 at  $z = 1.0$ ).
- 3) This implies the evolution of the UV-IR BLF is mainly due to the different evolution of the univariate LFs, and may not be controlled by the dependence structure.
- 4) The nonlinear structure of the BLF is naturally reproduced by the Gaussian copula.

The data from *Herschel*, *ALMA*, and *SPICA* data will improve the estimates drastically, and we expect to specify the full evolution of the UV-IR BLF in the Universe. We stress that the copula will be a useful tool for any other kind of bi-(multi-) variate statistical analysis.

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